

# **Chapter 29 An Experimental Study on Characteristics of Rolling in Head Waves for a Vessel with Nonlinear GZ-curve**

#### **Toru Katayama, Shugo Miyamoto, Hirotada Hashimoto and Yoshifumi Tai**

**Abstract** In this study, the characteristics of rolling in head waves for a vessel with strong nonlinear *GZ*-curve, which includes parametric rolling, are investigated. Rolling is measured for systematically changed wave length and height under the same forward speed, which is service speed in heavy weather. As a result, the range of  $T_e/T_\phi$  ( $T_e$  and  $T_\phi$  are encounter wave period and the roll natural period of the model) when oscillatory rolling occurs is wider than that of previous results by Taguchi et al. (Model Experiment on Parametric Rolling of a Post-Panamax Containership in Head Waves. Proceedings of the Ninth International Conference on Stability of Ships and Ocean Vehicles, [2006\)](#page-13-0), and the range spreads out wide area of  $T_e/T_\phi > 0.5$ . Especially, in the range of  $T_e/T_\phi > 0.5$ , oscillatory rolling is caused by large wave height and roll amplitude becomes larger than that at  $T_e/T = 0.5$ . It is supposed that the result is caused by change of roll natural period caused by nonlinear *GZ*-curve. In order to confirm it, numerical simulations are carried out for several variations of *GZ*curve. Additionally, roll measurements in irregular waves with Pierson-Moskowitz spectrum are also carried out.

**Keywords** Parametric rolling · Wide-breadth and shallow-draft Strongly nonlinear-GZ-curve

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### **29.1 Introduction**

Ballast water for vessels is significant to prevent their stability and propulsion performance losses, which are caused by their light-draught. However, the International Maritime Organization (IMO) adopted International Convention for the Control and Management of Ships' Ballast Water and Sediments in 2004, because of avoiding its environmental impact.

Then, in order to solve this problem, new hull forms are being designed, which use a small amount of ballast water or don't need ballast water. For example, one of new designed hull has shallow draft and wide breadth. Such hull may show different roll characteristics from the conventional vessels.

In the previous study (Tai et al. [2011\)](#page-14-0), for a shallow draft and wide breadth vessel, its roll characteristics in beam waves have been investigated experimentally. The results show that its roll natural period is changed by its roll amplitude and its roll resonance occurs in different wave period depending on different wave height.

In this study, the roll characteristics of the same model (Tai et al. [2011\)](#page-14-0) in head waves (including parametric rolling) are investigated. Roll measurements in regular head waves are carried out. From the results, it is found that large amplitude rolling occurs at a wider range of encounter frequencies than that of a conventional vessel, and it is supposed that it is caused by strong nonlinearity of its *GZ*-curves. Then, in order to be clear the reasons of the results, numerical simulations are carried out. Additionally, the roll characteristics in irregular waves with Pierson-Moskowitz spectrum are also investigated.

#### **29.2 Model and Measuerment**

Figure [29.1](#page-2-0) and Table [29.1](#page-2-1) show a body plan and principal particulars of a model. Figure [29.2](#page-2-2) shows the calculated *GZ*-curve of the model in real scale. In the same figure, the linear *GZ*-value of the model is also shown. The calculated *GZ*-curve shows strong nonlinear characteristics in the range of small roll angles. This is because bilge of the model is exposed above the water surface from the range of small roll angles.

Figure [29.3](#page-3-0) shows measured roll natural period of the model by a free decay test, which is carried out for various initial heel angles. This figure shows that the roll natural period becomes longer according to increase in roll amplitude. And this is caused by the strong nonlinearity of the *GZ*-curve shown in Fig. [29.2.](#page-2-2)

In order to investigate the occurrence of oscillatory rolling, towing tests in regular head waves are carried out at the towing tank of Osaka Prefecture University (length 70 m, breadth 3 m, depth 1.5 m). Figure [29.4](#page-3-1) shows a schematic view of the experiment. A model is attached to the towing carriage with two elastic ropes (Hashimoto et al. [2007\)](#page-14-1) . The model is towed at constant forward speed in head waves, and a small disturbance which is heel caused by hand is given. Roll motion is measured



<span id="page-2-0"></span>**Fig. 29.1** Body plan of the model



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by a gyroscope (CROSSBOW NAV440) and wave elevation is also measured by a servo type wave height meter attached to the towing carriage.

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#### **29.3 Roll Measurement in Regular Head Waves**

Table [29.2](#page-3-2) shows the experimental conditions. When the model is towed at  $F_n =$ 0.083 in head wave of  $T_w = 1.05$  s, the encounter wave period is half of the roll natural period ( $T_{\phi} = 1.88$  s) for small amplitude. This  $F_n$  is 7 kts in the real scale.

Figure [29.5](#page-4-0) shows examples of typical time histories of measured rolling in regular head waves. The upper figure shows the result that oscillatory rolling does not occur. The middle figure shows the result that oscillatory rolling occurs. The bottom figure shows the result that oscillatory rolling occurs after a small disturbance which is

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|                          | 0.083         |
|--------------------------|---------------|
| Wave period: $T_w$ [sec] | $0.99 - 1.55$ |
| $T_e/T_\phi$             | $0.42 - 0.71$ |
| Wave height: $H_w[m]$    | $0.01 - 0.05$ |

**Table 29.2** Experimental condition

 $T_e$ : encounter wave period,  $T_\phi$ : roll natural period

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heel caused by hand, this phenomenon is investigated as a bi-stability mechanism by Spyrou et al. [\(2008\)](#page-14-2).

Figure [29.6](#page-5-0) shows measured roll amplitudes in regular head waves. In this figure, its horizontal axis is the ratio of encounter wave period to the roll natural period for small amplitude. At  $T_e/T_\phi = 0.5$ , oscillatory rolling occurs for any wave heights. For small wave height  $H_w = 0.01$  m, oscillatory rolling does not occur except at  $T_e/T_\phi$  $= 0.5.$ 

Figure [29.7](#page-5-1) shows the measured roll amplitude of a conventional post-panamax container vessel by Taguchi et al. [\(2006\)](#page-13-0) in order to compare with Fig. [29.6.](#page-5-0) From the comparisons, it is noted that the range of  $T_e/T_\phi$  where oscillatory rolling occurs in Fig. [29.6](#page-5-0) is wider than that of Fig. [29.7](#page-5-1) even if its amplitude is smaller, and the range of Fig. [29.6](#page-5-0) spreads wider area of  $T_e/T_\phi > 0.6$ . Especially, at the range over  $T_e/T_\phi = 0.5$ . in Fig. [29.6,](#page-5-0) oscillatory rolling is caused by large wave height, and its amplitude becomes larger than that at  $T_e/T_\phi = 0.5$ . It is supposed that the spread of occurrence of large amplitude rolling to wider area of  $T_e/T_\phi > 0.5$  in Fig. [29.6](#page-5-0) is caused by the strong nonlinearity of *GZ*-curves.

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<span id="page-5-2"></span>Table [29.3](#page-5-2) shows the experimental condition for the roll measurement in regular head waves with constant wave length ( $\lambda/L_{pp} = 1$ ). When model is towed at  $F_n =$ 0.077, its encounter wave period is half of the roll natural period for small amplitude  $(T_\phi = 1.88 \text{ s})$ . According to decrease in forward speed, its encounter period becomes long.

Figure [29.8](#page-6-0) shows the measured roll amplitude. In this figure, roll amplitude increases according to decrease in forward speed, even if  $T_e/T_\phi$  becomes larger than 0.5.

In order to make its reason clear, the characteristics of roll damping is investigated by using Ikeda's method, and roll amplitude is estimated by using a simplified estimation for amplitude of parametric rolling (Katayama et al. [2009\)](#page-14-3) .

Figure [29.9](#page-6-1) shows the estimated roll damping. The roll damping decreases according to decrease in forward speed. It is caused by decrease in the lift component of roll damping which is caused by forward speed.

Figure [29.10](#page-6-2) shows the estimated roll amplitude by Eq. [\(29.1\)](#page-6-3) (Katayama et al. [2009\)](#page-14-3).

<span id="page-6-1"></span><span id="page-6-0"></span>

<span id="page-6-3"></span><span id="page-6-2"></span>The measured results are also shown in the same figure. Both results show the same tendency for change of forward speed.

From these results, it is understood that decrease in forward speed causes decrease in roll damping, and roll amplitude becomes larger. On the other hand, roll natural period becomes long according to increase in roll amplitude as shown in Fig. [29.3.](#page-3-0) Then rolling resonates with increased wave encounter period by decrease in forward speed, and parametric rolling occurs.

#### **29.4 Roll Measuring in Irregular Head Waves**

Equation [\(29.2\)](#page-7-0) is one of Pierson-Moskowitz spectrum.

<span id="page-7-0"></span>
$$
S(\omega) = \frac{8.1 \times 10^{-3} \times g^2}{\omega^2} \exp\left(-\frac{3.11/H_{1/3}^2}{\omega^4}\right)
$$
 (29.2)

To make irregular waves, the spectrum is divided into 1000 equally in  $\omega =$ 1–15 rad/sec and a sine wave of each frequency component is superposed. In addition, the phase difference of each frequency component is given as random numbers. The number of measurements is decided by the convergence of variance of rolling (Appendix).

In this study, roll measurements are carried out for three significant wave heights, and Fig. [29.11](#page-7-1) shows the power spectral densities obtained from the measured wave elevations.

Figure [29.12](#page-8-0) shows the probability distributions of the measured roll amplitudes. In this measurement, roll amplitude is not over 10 deg at significant wave height  $=$ 0.04 m. The results show that the roll amplitude increases with increase in significant wave height.

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<span id="page-8-1"></span><span id="page-8-0"></span>

Figure [29.13](#page-8-1) shows time histories of measured wave height and rolling at  $H_w$  = 0.06 m. In this figure, the ratios of mean encounter wave period to the roll natural period for small amplitude  $(T_e/T_\phi)$  at  $t_1$  and  $t_2$  are also shown. The mean wave heights at  $t_1$  and  $t_2$  are almost same. From this figure, it is found that oscillatory rolling does not occur at  $t_1(T_e/T_\phi < 0.5)$ , and oscillatory rolling occurs at  $t_2(T_e/T_\phi$  $> 0.5$ ).

In the Pierson-Moskowitz spectrum expressed by Eq. [\(29.2\)](#page-7-0), increase in significant wave height causes increase in mean wave period, therefore the occurrence probability of a wave group containing frequency component of  $T_e/T_\phi > 0.5$  mainly also becomes higher. On the other hand, the roll natural period becomes longer according to increase in roll amplitude. Therefore, rolling may resonate the wave group containing frequency component of  $T_e/T_\phi > 0.5$  mainly.

## **29.5 Roll Measuring in Head Waves of Two Different Wave Period**

In the roll measurement in irregular head waves, when mean wave period of group waves is  $T_e/T_\phi > 0.5$ , oscillatory rolling occurs. In order to investigate its reasons, the roll measurements are carried out in the waves which are generated by the superposition of two sine waves having different periods. One of two different periods is half ( $T_e/T_\phi = 0.5$ ) of the roll natural period for small amplitude. Moreover, the wave heights of two sine waves are different. (In the following, the component with higher wave height is called Main-wave.) Experimental conditions are shown in Table [29.4](#page-9-0) with the results.

In Table [29.4,](#page-9-0) measured roll amplitude is shown. The results show that oscillatory rolling occurs when period of Main-wave is  $T_e/T_\phi = 0.50{\text -}0.61$ . As the results, it is confirmed that rolling can resonate with a wave with  $T_e/T_\phi = 0.50{\text -}0.61$  and it is supposed that oscillatory rolling may occur in a frame of wide line.

#### **29.6 Numerical Simulation**

By utilizing the numerical simulation model (Hashimoto and Umeda [2010\)](#page-13-1) , the effects of nonlinearity of *GZ*-curve on the occurrence of parametric rolling are investigated. In the simulation model, a 3DOF sway-heave-roll motion is solved in time-domain based on a nonlinear strip method in which dynamic components, i.e. radiation and diffraction forces, are calculated for an asymmetric submerged hull due to roll. Hydrodynamic forces for the heave and diffraction modes are determined

<span id="page-9-0"></span>

|   | $T_e/T_\phi$ ( $H_w = 0.01$ m) |      |          |          |      |      |          |          |  |
|---|--------------------------------|------|----------|----------|------|------|----------|----------|--|
| $\frac{T_e/T_{\phi}}{(H_w=)}$<br>$0.04 \text{ m}$ ) |                                | 0.39 | 0.45     | 0.50     | 0.55 | 0.61 | 0.66     | 0.71     |  |
|   | 0.39                           |      |          | $\times$ |      |      |          |          |  |
|   | 0.45                           |      | $\times$ | $\times$ |      |      |          |          |  |
|   | 0.50                           | 10.4 | 14.1     |          | 18.4 | 15.4 | 18.3     | 16.8     |  |
|   | 0.55                           |      |          | 8.3      |      |      |          |          |  |
|   | 0.61                           |      |          | 16.0     |      |      |          |          |  |
|   | 0.66                           |      |          | $\times$ |      |      | $\times$ |          |  |
|   | 0.71                           |      |          | $\times$ |      |      |          | $\times$ |  |

**Table 29.4** Measred results

 $_{\phi}$  [deg]: oscillatory rolling occurs

×: oscillatory rolling does not occur

Colored: oscillatory rolling occurs in regular waves

at the encounter frequency and those for the sway and roll modes are done at the half the encounter frequency.

In order to control the nonlinearity of *GZ*-curve, the breadth-draft ratio (*B/d*) of original model is changed under the constant displacement. Figures [29.14](#page-10-0) and [29.15](#page-10-1) show modified body plans and *GZ*-curves. *GM* for each modified model is adjusted to the same *GM* as the original model. Moreover, in order to make sure the characteristics of Δ*GM* of the original and modified models, the estimated Δ*GM* by the simplified method (Umeda et al. [2011\)](#page-14-4) are shown in Fig. [29.16.](#page-11-0)

Figure [29.17](#page-11-1) shows the numerical results for the original model. The calculated results show large amplitude parametric rolling occurs in the wide range over  $T_e/T_\phi$ = 0.5 and the numerical result shows a good agreement with the measured results.

Figures [29.18,](#page-11-2) [29.19](#page-11-3) and [29.20](#page-12-0) shows the calculated results for the modified models. These results show the roll amplitudes in the range over  $T_e/T_\phi = 0.5$  become larger according to the increase in nonlinearity of *GZ*-curve.

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<span id="page-11-1"></span><span id="page-11-0"></span>

# <span id="page-11-3"></span><span id="page-11-2"></span>**29.7 Conclusions**

In this study, the roll characteristics of a wide breadth and shallow draft vessel in head waves are investigated. And the following conclusions are obtained.

<span id="page-12-0"></span>

- In regular head waves, a oscillatory rolling occurs at wider range of encounter frequency than that of a conventional vessel. Especially, in the range of  $T_e/T_\phi$ > 0.5, a oscillatory rolling is caused by large wave height and roll amplitude becomes large.
- Numerical simulations demonstrate that the large amplitude parametric rolling occurs in the wide range over  $T_e/T_\phi = 0.5$  because of the strong nonlinearity of *GZ*-curves.
- There is the possibility that slowdown increases roll amplitude, when parametric rolling occurs at  $\lambda/L_{pp} = 1.0$ .
- It is confirmed by simplified estimation that this phenomenon is caused by the following scenario. Slowdown decreases roll damping, and roll amplitude becomes large. On the other hand, roll natural period becomes longer with increase in roll amplitude. Therefore, rolling resonates with increased wave encounter period by slowdown.
- In irregular head waves with a Pierson-Moskowitz spectrum, a oscillatory rolling is caused by resonating with the group waves which contain the frequency component of  $T_e/T_\phi > 0.5$  mainly.

#### **Appendix**

#### *The Number of Times of Test in Irregular Head Waves*

From wave height, pitch motion and roll motion, each variance and ensemble mean by Eqs. [\(29.3\)](#page-13-2), [\(29.4\)](#page-13-3). Figure [29.21](#page-13-4) shows calculated result. In these figure, a vertical axis is a ratio ensemble mean to standard deviation. Upper and middle figure uses 20 samples and bottom figure uses  $30$  samples. (1 sample  $=$  40 s) From upper and middle figure, it is found that wave height and pitch motion convergence with 5 samples. On the other hand, roll motion does not converge until 15 samples. So, this study uses 20 samples.

<span id="page-13-4"></span>

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