# Chapter 18

## **Directed Modification of Reaction Centers from Purple Bacteria**

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### **Summary**

 Reaction centers from purple bacteria form a superb test system for the manipulation of electron transfer parameters. The wealth of cofactors and electron transfer reactions provides opportunities for directed modification of specifi c properties. In particular, the energies of each cofactor can be selectively changed by mutations of neighboring amino acid residues. The starting point for the initial electron transfer, the bacteriochlorophyll dimer, has proven to be exceptionally malleable, allowing large changes in energetics and rates. Most of the other cofactors can be exchanged or eliminated entirely, permitting considerable alteration of pathways. By orchestrating multiple changes in the reaction center, the light-initiated electron transfer pathway can be directed towards alternate ends, for example down the B branch of cofactors rather than the naturally preferred A branch. Extensive modeling of features of electron transfer such as the energetics, the coupling, and the protein dynamics has been corroborated by observed changes in the characteristics of the reactions after modification of the cofactor properties. For example, the maximum rates for several electron transfer reactions, determined by application of Marcus theory to the rates of reactions in a range of mutants, show a correlation with the

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distance between the cofactors. Other measurements revealing the intimate interaction of the protein and cofactors show that protein motion controls the rate of the initial electron transfer. Thus the reaction center provides a natural and modifiable template for understanding the factors governing electron transfer.

#### **I. Introduction**

 What happens when the reaction center is excited by light? The light energy is converted into chemical energy through a series of electron and proton transfer reactions involving the cofactors of the reaction center. These reactions are able to proceed with essentially every light photon producing useful reactions, corresponding to a quantum efficiency of nearly 100%. The balancing act of capturing light energy while not destroying the molecules involved or producing unfavorable side reactions is achieved by fine tuning the properties of the cofactors through interactions with the protein in which they are embedded. The protein can influence several aspects of electron transfer identified by theoretical treatments, including the energetics, the coupling, and the protein dynamics. Although it is difficult to separate the contributions of each of these components, experiments have probed their effects, notably by altering the light-induced reactions in specific ways. This chapter will review the partnership between the cofactors and the protein scaffold as it relates to the parameters of electron transfer models, focusing on examples where changes in the properties have measurable effects on electron transfer. The examples will be primarily of reaction centers from *Rhodobacter* (*Rba.*) *sphaeroides* and*Rba. capsulatus*, the two commonly studied reaction center systems of purple bacteria.

#### **II. Properties of the Cofactors**

#### A. Identity, Substitution, and Removal

The reaction center from *Rba. sphaeroides* and *Rba. capsulatus* is an integral membrane protein complex composed of three proteins: the L, M, and H subunits. The L and M subunits form the core of the protein and are largely composed of five transmembrane helices that are structurally related to each other by an approximate two-fold symmetry axis. The H subunit is more peripheral, containing one transmembrane helix and a large cytoplasmic domain. Embedded in the middle of the L and M subunits are ten cofactors, all of which can participate in some manner in energy or electron transfer. The ten cofactors in the reaction center are: two bacteriochlorophyll *a* molecules that form a dimer (P), two bacteriochlorophyll *a* monomers ( $B_A$  and  $B_B$ ), two bacteriopheophytin *a* molecules ( $H_A$  and  $H_B$ ), two ubiquinone molecules  $(Q_A$  and  $Q_B$ ), a carotenoid molecule, and an iron (Fig. 1). These cofactors are arranged into two branches, identified as the A and B branches, which are related by the same two-fold symmetry as found for the L and M subunits.

Although the cofactors are normally expressed with a well-defined composition, some of these cofactors can be substituted by molecules in the same class, for example bacteriochlorophylls for bacteriopheophytins. The monomer bacteriochlorophylls and bacteriopheophytins, the quinones, the carotenoid, and the iron can be biochemically removed and replaced. Mutagenesis can also result in biosynthetic substitutions, primarily by replacement of the amino acid residues coordinating the cofactors. For example, when the ligand to one of the central Mg atoms of P is changed, a bacteriopheophytin is incorporated rather than bacteriochlorophyll in a mutant that has been termed a heterodimer (His L173 to Leu and His M202 to Leu in *Rba. sphaeroides*, L173 and M200 in *Rba. capsulatus*) (Bylina and Youvan, 1988; Kirmaier et al., 1988; McDowell et al., 1991; Allen et al., 1996; van Brederode et al., 1999; King et al., 2001). Both halves of the dimer can be individually replaced this way, although the double bacteriopheophytin dimer appears to be unstable. Similarly the binding site for the B-side bacteriochlorophyll monomer is found to contain a bacteriopheophytin when the residue forming its Mg ligand is changed ( $\phi$  mutant, His M182 to Leu in *Rba. sphaeroides*), although analogous mutations on the A side do not appear to have the same substitution effect (Katilius et al., 1999, 2004). Conversely, the bacteriopheophytin on the A side can be converted to bacteriochlorophyll by introduction of a residue to act as a ligand  $(\beta \text{ mutant}, \text{Leu M214})$ 

*Abbreviations:*  $B_A$  – bacteriochlorophyll monomer on A branch of cofactors;  $B_B$  – bacteriochlorophyll monomer on B branch of cofactors;  $H_A$  – bacteriopheophytin on A branch of cofactors;  $H_B$  – bacteriopheophytin on B branch of cofactors; P – bacteriochlorophyll dimer;  $Q_A$  – quinone on A branch of cofactors; QB – quinone on B branch of cofactors; *Rba. – Rhodobacter*



*Fig. 1.* Three-dimensional structure of the cofactors of the reaction center from *Rba. sphaeroides* R-26. Shown are the bacteriochlorophyll *a* dimer (P) (shaded dark), the two bacteriochlorophyll *a* monomers ( $B_A$  and  $B_B$ ), the two bacteriopheophytin *a* molecules  $(H_A \text{ and } H_B)$ , the two ubiquinone molecules  $(Q_A \text{ and } Q_B)$ , and the iron (Fe). Although the carotenoid is present in wild type, it is not present in the R-26 strain. The view is perpendicular to the approximate two-fold symmetry axis that passes from P to Fe in the plane of the paper.

to His in *Rba. sphaeroides*, M212 in *Rba. capsulatus*) (Kirmaier et al., 1991). The addition of a histidine near the bacteriopheophytin on the B branch also results in incorporation of a bacteriochlorophyll (Leu L185 to His in *Rba. sphaeroides*) (Watson et al., 2005). Changing iron ligands results in a loss of metal specificity, and in one case a significant amount of zinc is incorporated (His M266 to Cys in *Rba. sphaeroides*) (Williams et al., 2007).

In addition to altering the cofactor composition, some of the cofactors can be removed or their incorporation can be blocked. Many studies of reaction centers have been performed on a carotenoid-less strain of *Rba. sphaeroides*, identified as R-26, which shows properties essentially identical to the carotenoid-containing wild type except for the loss of the ability to trap excess energy. The quinones can be taken out by exposing the reaction centers to a detergent treatment, which initially results in a decrease in  $Q_B$  followed by loss of  $Q_A$ . Biosynthetic incorporation of  $Q_A$  can be blocked by substitution of amino acid residues forming the binding pocket,

for example by removing the tryptophan in van der Waals contact with  $Q<sub>A</sub>$  (M252 in *Rba. sphaeroides* and M250 in *Rba. capsulatus*), or adding a tryptophan in place of a smaller residue (Ala M260 to Trp in *Rba. sphaeroides*) (Breton et al., 2004). Certain mutations near P result in reaction centers that lack a functional P (Val L157 to Arg, His L153 to Glu, Leu, Gln, or Tyr, His L173 to Gly, and His M202 to Gly in *Rba. sphaeroides*) (Jackson et al., 1997; Moore and Boxer, 1998; Katilius et al., 2004). The loss of  $H_A$  is one outcome of the large-scale alterations of the  $D_{LL}$ mutant, in which the D transmembrane sequence of the M subunit is replaced with the symmetry-related segment of the L subunit (M192 to M217 replaced with L165 to L190 in *Rba. capsulatus*) (Robles et al., 1990). Similarly, the B-branch bacteriopheophytin is not required for assembly of the reaction center as shown by a mutant with the change of an alanine that is adjacent to  $H<sub>B</sub>$  to tryptophan (M149) in *Rba. sphaeroides*)(Watson et al., 2005). A loss of bacteriochlorophyll (presumably in the dimer) was also reported to be due to structural and electrostatic changes in a residue located between  $B_B$  and P (Ile L177 to His in *Rba. sphaeroides*) (Khatypov et al., 2005). See Chapter 16, Jones, for a summary of the effects of exclusion and replacement of reaction center cofactors. The ability to alter the cofactor composition provides the opportunity to manipulate the electron transfer reactions as discussed below.

#### B. Optical Spectra

One of the most accessible properties of the cofactors is the absorption spectrum (Fig. 2). The tetrapyrrole pigments (P,  $B_A$ ,  $B_B$ ,  $H_A$ , and  $H_B$ ) have absorption peaks in the near-infrared region, the visible region and the UV region, and the quinones have an unresolved band in the visible region. The bacteriopheophytins, monomer bacteriochlorophylls, and dimer bacteriochlorophylls can be distinguished from each other in the near-infrared peaks at 760 nm, 800 nm, and 865 nm, respectively. In the visible region, the 540 nm peak arises from the bacteriopheophytins, and the 590 nm peak is from all four bacteriochlorophylls. The A and B branch pigments of the same type overlap, except at low temperature where the broad peak in the 540 nm region of the spectrum is resolved into two peaks at 533 nm and 546 nm associated with  $H_B$  and  $H_A$ , respectively. The contributions of the tetrapyrrole pigments in the Soret region have been delineated, with H contributing primarily on the



*Fig. 2.* Absorption spectrum of reaction centers from wild-type *Rba. sphaeroides*. The primary contributions of the bacteriochlorophyll dimer (P), bacteriochlorophyll monomers (B), and bacteriopheophytin monomers (H) to each of the absorption bands are identified. The Soret band arises from all the tetrapyrrole cofactors.

blue side and P absorbing on the red side (Wang et al., 2006). The carotenoid has a major but poorly resolved contribution to the absorption near 500 nm, where absorption of the other cofactors is weak, and can transfer excitation energy to the bacteriochlorophylls, performing a light-harvesting function in addition to photoprotection (Lin et al., 2003).

The absorption peaks in the near-infrared region are shifted from the solution spectra, notably in the shift of the bacteriochlorophyll bands to longer wavelengths, presumably because of their protein environment. However, specific alteration of the spectrum generally has not been amenable to mutagenesis. Modeling of the excited states has also proven difficult (Dahlbom and Reimers, 2005). However in certain instances, shifts in the spectrum are attributed to changing particular residues. A notable example is the change in the visible region peak of the bacteriopheophytins due to changing a hydrogen bond to the keto group of  $H<sub>A</sub>$ , which established the assignment of these optical bands to the individual bacteriopheophytins (Bylina et al., 1988). However, changing a hydrogen bond is not typically correlated with a shift in the peak of the tetrapyrrole pigments. Another relatively malleable absorption peak is that of the dimer. It can shift up to approximately 15 nm to shorter wavelengths as a result of mutations, mostly mutations in which a hydrogen bond to the acetyl group is changed, precipitating a rotation of this side group. Major changes in the absorption spectra also occur when substitutions

of the cofactors are introduced. For example, the nearinfrared absorption peak of the dimer is significantly different in the heterodimer mutant. Likewise, shifts in the both the visible and near-infrared regions of the spectra are observed in mutants with alterations of ligands to the monomer bacteriochlorophylls that result in pigment changes (Katilius et al., 1999, 2004). Thus, the optical spectrum is a sensitive indicator of the effects of certain types of modifications to the reaction center.

The absorption bands in the near-infrared region arise from transitions from the ground state to the first excited state and so are markers of the excited state energy, indicating the maximum amount of energy that can be captured. For the primary donor of *Rba. sphaeroides*, the absorption peak at 865 nm corresponds to an energy difference of 1.4 eV. The properties of the reaction center can also be characterized from measurement of the spontaneous and stimulated emission of the excited state of the dimer, centered near 915 nm. Because the absorption and emission bands change as the cofactors undergo excitation, oxidation, and reduction, transient optical spectroscopy is one of the major techniques utilized to follow the light-induced transfer of electrons in the reaction center.

#### C. Oxidation/Reduction Midpoint Potentials

The oxidation/reduction midpoint potentials of the reaction center cofactors are critical properties for their function as electron transfer components. The midpoint potential of the dimer, at approximately 500 mV, is the only one easily measured directly (Fig. 3). The midpoint potentials of the other cofactors can only be inferred. A change in the chemical nature of a cofactor has a direct effect on its midpoint potential. For example, incorporation of a bacteriopheophytin in the heterodimer mutant increases the potential by approximately 130 mV due to the intrinsically higher potential of bacteriopheophytin (Allen et al., 1996). In addition, the energies of the electronic states are sensitive to the environment so protein interactions with the dimer, including hydrogen bonds and electrostatic forces from charged residues, can affect the dimer midpoint potential and be modulated by mutagenesis. The midpoint potentials of other tetrapyrroles in the reaction center can presumably also be changed by similar modifications, although the evidence is based upon alterations of the electron transfer rates rather than direct measurements.



*Fig. 3.* (Left) The near-infrared region of the absorption spectrum showing a systematic decrease in the band associated with the bacteriochlorophyll dimer as the ambient potential is increased. The absorbance of this band is used to determine the fraction of  $P^+$  at any given potential compared to the total amount of P. (Right) Fits of the dependence of the fraction of  $P^+$  on the ambient potential are used to determine the oxidation-reduction midpoint potential,  $E_m$ . A shift in the midpoint potential is observed for a number of different mutants (Allen and Williams, 1995).

Whereas some modifications lead to large changes, analysis of various mutations indicates that most result in small  $(<50 \text{ mV})$  increases in the  $P/P^+$  midpoint potential (Spiedel et al., 2002). The predominance of increases in the midpoint potential of P, rather than a mix of increases and decreases, is probably because the midpoint potential is poised by the protein interactions at a minimum value, and perturbations primarily disrupt this poise.

The effect of hydrogen bonds to the conjugated carbonyl molecules on the oxidation/reduction midpoint potential of the dimer is well documented (Allen and Williams, 1995, 2006). An additional hydrogen bond raises the potential by 60 to 120 mV while loss of the hydrogen bond present in the wild type decreases the potential by approximately 80 mV (Stocker et al., 1992; Williams et al., 1992; Murchison et al., 1993). Multiple systematic changes in the hydrogen bonding pattern produce a wide range of midpoint potentials (Lin et al., 1994a) and provide the opportunity to investigate the effect of altered energetics on the properties of the reaction center as discussed below. Hydrogen bonds have been introduced to the monomer bacteriochlorophylls, but with little effect on the primary photochemistry (Chen et al., 2004). A water molecule is found in the wild type between P and  $B_A$  in hydrogen bonding position to the keto carbonyl of  $B<sub>A</sub>$ . This water molecule can be displaced by mutations (at M203 in *Rba. sphaeroides*), resulting in alteration in the electron transfer characteristics, including a slowing of the primary electron transfer rates, suggesting that  $B_A$  is more difficult to reduce, although other factors, such as an increase in the reorganization energy (see below), would also result in a slower rate (Potter et al., 2005; Yakovlev et al., 2005). Hydrogen bonds to bacteriopheophytins are also probably influencing their midpoint potentials. On the A side, a hydrogen bond between the  $13<sup>1</sup>$  keto of  $H_A$  and a glutamic acid residue (L104 in *Rba*. *sphaeroides* and *Rba. capsulatus*) is naturally occurring, and the equivalent situation on the B side can be achieved by mutagenesis (Val M131 to Asp in *Rba. capsulatus*, M133 in *Rba. sphaeroides*) (Bylina et al., 1988; Müh et al., 1998; Kirmaier et al., 2002a). The presence of the hydrogen bonds most likely makes the bacteriopheophytins easier to reduce, consistent with the observed changes in electron transfer rates in the mutants.

The potential of each cofactor is also shaped by electrostatic interactions with charged and polar amino acid residues. Insertion or removal of ionizable residues at several different locations approximately 10 to 15 Å from P (L135, L155, L164, L170, L247 and M199 in *Rba. sphaeroides*) leads to a midpoint potential decrease up to 60 mV due to a negative charge and an increase up to 50 mV due to a positive charge (Williams et al., 2001; Johnson and Parson, 2002; Johnson et al., 2002). The effect of the charges on the potential is relatively modest because of screening of the charge-charge interactions by the surrounding protein. The energies of  $B_A$  and  $B_B$  can be significantly changed by altering the polarity of residues near these tetrapyrroles, including replacing a conserved tyrosine residue with phenylalanine (M210 in *Rba. sphaeroides*, M208 in *Rba. capsulatus*), and changing a conserved phenylalanine to tyrosine (L181 in *Rba. sphaeroides* and *Rba. capsulatus*), as measured by changes in the electron transfer rates (Finkele et al., 1990; Nagarajan et al., 1990; Jia et al., 1993) and Stark spectra (Treynor et al., 2004). Theoretical calculations suggest that electrostatic fields near  $B_A$ and  $B_B$  are influenced by the polarity of these residues, and consequently mutations result in energetic shifts for the oxidized states (Alden et al., 1996; Gunner et al., 1996). In general, the sensitivity of the cofactors to specific protein interactions provides a means to test the role of the energies of the charge-separated states in determining the electronic structure of the cofactors and the rates of electron transfer as presented in the subsequent sections.

#### D. Modeling the Electronic Structure of the Bacteriochlorophyll Dimer

The close overlap of the two tetrapyrroles in P results in a sharing of the electron orbitals and hence changes in the properties of P compared to bacteriochlorophyll monomers. For example, one effect of the dimerization is a shift in the absorption band to a longer wavelength than is observed for the bacteriochlorophyll monomers. In a simple Hückel molecular orbital model, the two conjugated molecules can be considered to be coupled together, resulting in the electrons being distributed over the two bacteriochlorophylls in molecular orbitals that

are split by an energy 2β (Fig. 4). An additional term represents the difference in the energies of the two sides arising from the inhomogeneous nature of the protein surrounding P. When P is oxidized, the higher molecular orbital loses an electron, leaving one unpaired electron. In the Hückel model, the predominant contribution to the higher molecular orbital is from the side with the higher energy (Plato et al., 1992). In the wild-type reaction center, measurement of the unpaired electron spin densities using the magnetic resonance technique called electron nuclear double resonance yields a 2:1 ratio for the spin density on the L side bacteriochlorophyll compared to the M side. The energy difference between the two sides can be manipulated by introduction of hydrogen bonds to P, resulting in systematic changes in the molecular energies and hence the ratio of the spin densities (Fig. 4) (Artz et al., 1997; Müh et al., 2002). In general, a hydrogen bond between the side chain of an amino acid residue and the M side of P stabilizes the energy of that side and increases the energy difference of the molecular orbitals, making the spin ratio more asymmetric. A hydrogen bond to the L side of P stabilizes that side and hence decreases the energy difference for the molecular orbitals and produces a more symmetrical spin distribution. In both cases, the additional hydrogen bond lowers the energy of the highest molecular orbital, which is coupled to the midpoint potential, making P more difficult to oxidize. More extensive modeling involving additional factors, such as the contribution of vibrational states, provides estimates for the positions of bands seen in the infrared region of the optical spectrum (Müh et



*Fig. 4.* A two-orbital Hückel model of the bacteriochlorophyll dimer. In wild type, the energy splitting of the molecular orbitals of P is determined by the energy difference between the two sides of P and the coupling. Mutations that introduce a hydrogen bond to the M side of P stabilize the energy of the M side bacteriochlorophyll, resulting in a more asymmetrical dimer with a larger energy difference for the molecular orbitals and a larger midpoint potential (*E*m). Introducing a hydrogen bond to the L side results in a more symmetrical dimer while still increasing *E*m.

al., 2002; Reimers and Hush, 2004; Kanchanawong et al., 2006).

In comparison to P in the bacterial reaction center, the electronic structure of the primary electron donor of Photosystem I, P700, is very asymmetrical with a ratio of at least 3:1 in spin density distribution over the two tetrapyrroles (Webber and Lubitz, 2001). Two factors contribute to this asymmetry. The primary donor has a heterodimeric nature as the Bside tetrapyrrole is a chlorophyll *a*´, which contains an epimeric configuration at the  $13<sup>2</sup>$  position, rather than a chlorophyll *a* as found on the A side. Whereas both sides are axially coordinated by histidines, only the chlorophyll *a*´ has a hydrogen bond (between Thr 739 of the PsaA subunit and the  $13<sup>1</sup>$  position). The presence of this hydrogen bond likely stabilizes the chlorophyll *a*´ side of P700 resulting in the chlorophyll *a* side having a much higher energy and hence a much larger unpaired spin density. Mutation of Thr 739 to Ala results in a 60 mV decrease in the oxidation/reduction potential while a 30 mV decrease is observed for a Val mutation (Witt et al., 2002; Li et al., 2004). Associated with the loss of the hydrogen bond is a small redistribution of the electron spin density from the B side to the A side of P700 as expected based upon the theoretical model. Thus, the understanding developed from the bacterial system appears to be applicable to other photosystems.

#### **III. Electron Transfer Concepts**

According to Marcus theory, the rate, *k*, of electron transfer between a donor and acceptor can be written in terms of the temperature, *T*, the effective electronic matrix element, or coupling, *V*, and energetics according to:

$$
k = \sqrt{\frac{4\pi^3}{h^2 \lambda RT}} V^2 e^{-\left(\frac{(\Delta G^\circ + \lambda)^2}{4\lambda RT}\right)}
$$
(1)

with two contributions to the energetics: the free energy difference between the final and initial states, ∆*G*°, and the reorganization energy, λ, which represents the rearrangement energy associated with the charge transfer (Marcus and Sutin, 1985). The exponential dependence of the rate on the energy terms is due to the activation energy associated with the process. Electrons can flow over long distances in biological systems because of a careful balance, mediated by the protein, involving the free energy difference, the reorganization energy and the coupling between distant redox cofactors. The factors that influence the energetics and coupling in protein complexes are considered below.

#### A. Energetics

The electron transfer rate has an exponential dependence upon the free energy difference and reorganization energy (Eq. 1). The free energy differences between various states formed during electron transfer steps in wild type range from relatively small values of approximately 200 meV for the initial electron transfer to a value of approximately 500 meV for charge recombination from the primary and secondary quinones. Mutants with altered  $P/P^+$  midpoint potentials have corresponding changes in the free energy differences for the electron transfer steps. For example, a higher  $P/P^+$  midpoint potential increases the free energy difference between the  $PQ_A$  ground state and the  $P^+Q^-$  state. By measuring the rates for different mutants, the dependence of the electron transfer rates on the free energy difference can be experimentally determined. The dependence of the electron transfer rate typically is shown as the logarithm of the rate versus the free energy difference, yielding a parabolic curve (Fig. 5). The value of the free energy difference at the peak of the curve is equal to the reorganization energy, because for that value the exponential term has a maximal value of one. The rate at the peak of the parabola is proportional to the coupling for the reaction.

In principle, the energies of the cofactors, including their charged states, can be established by theoretical calculations based upon the three-dimensional structure of the protein. However, for large pigmentprotein complexes, these efforts remain a challenge due to the large number of interactions and the uncertainties in the atomic positions often found in structural models that are limited by the resolution of the X-ray diffraction data. Electrostatic potentials calculated for the reaction center from *Blastochloris viridis* show a strong asymmetry in the electrostatic potential for the two branches, with the A branch being substantially more positive than the B branch, significantly favoring electron transfer along the A branch (Gunner et al., 1996). The oxidation/reduction midpoint potential for P680 of Photosystem II is calculated to have a significantly higher potential than P in bacterial reaction centers, in agreement with



*Fig. 5.* (Left) The dependence of the electron transfer rate on the free energy difference for four reactions: the initial P\* decay, electron transfer from bound cytochrome  $c_2$  to  $P^+$ , and charge recombinations from the primary and secondary quinones. The lines are fits according to the Marcus relationship (Eq. 1). The data points and fits are from previous publications (Lin et al., 1994b; Allen et al., 1998; Haffa et al., 2002). (Right) The maximum rate from each fi t is proportional to the coupling for the reaction and yields a rate that is predicted to be exponentially related to the separation distance between the electron donor and acceptor. For the P\* decay, the maximum rate is plotted for transfer to two possible acceptors,  $B<sub>A</sub>$  and  $H<sub>A</sub>$ . The fit shown is from Dutton and coworkers (Noy et al., 2006).

the experimental data (Ishikita et al., 2006). These calculations also suggest that the specific orientation of the 132 ester of bacteriochlorophylls has a strong influence on the potential (Ishikita et al., 2005). Calculations making use of dynamics and alternate values for the dielectric constant result in different values for the potentials but are still in agreement with favorable transfer along the A branch (Parson et al., 1990; Gehlen et al., 1994).

Experimentally, the strength of the electrostatic fields can be established using a technique known as electroabsorption or Stark spectroscopy, in which changes in the optical spectrum due to the application of an electric field across a sample are measured. These optical changes can be interpreted in terms of the interactions of the applied electric field with the dipole moments of the cofactors. For the bacterial reaction center, such measurements lead to estimates of a strong field of  $10<sup>6</sup>$  V/cm near the bacteriochlorophyll dimer (Middendorf et al., 1993). The electrostatic field around P changes by approximately  $10\%$ when comparing reaction centers with and without the carotenoid (Yanagi et al., 2005). Because the electron transfer rates are insensitive to the presence or absence of the carotenoid, the electrostatic field is not a major determinant for the directionality along the A branch.

Replacement of one of the coordinating histidine residues of P with leucine produces the heterodimer, while mutants with the histidine replaced with glycine are very similar to wild type in their properties, presumably due to the incorporation of water as a ligand in place of the histidine (M202 in *Rba. sphaeroides*) (Goldsmith et al., 1996). Likewise, replacement of the histidine coordinating  $B_A$  with a small amino acid residue (His L153 to Ser or Gly in *Rba. sphaeroides,*  Thr or Ser in *Rba. capsulatus,* and Cys in *Blastochloris viridis*) results in mutants with properties that are very similar to those of wild type, while mutation to leucine alters the pigment composition (Bylina et al., 1990; Arlt et al., 1996; Katilius et al., 2004). Thus, these bacteriochlorophylls are tolerant of coordination changes, with the energetics of P and  $B<sub>A</sub>$  being only weakly dependent upon the nature of the coordination.

Protein interactions such as hydrogen bonds to the intermediate electron acceptors also influence the energetics of electron transfer. Removal of the hydrogen bond between a glutamic acid residue (L104 in *Rba. sphaeroides* and *Rba. capsulatus*) and the keto group at the  $13<sup>1</sup>$  position of  $H_A$  results in a small decrease in the forward electron transfer rate from P\* (Bylina et al., 1988), consistent with a decrease in the free energy difference due to a change in the

energy of  $H_{A}^-$ . In Photosystem I, loss of the hydrogen bond between a tyrosine residue and either of the chlorophyll monomers located in positions analogous to  $H_A$  and  $H_B$  (Tyr 696 of the PsaA subunit near the 13<sup>1</sup> position of ec3<sub>A</sub> or Tyr 676 of the PsaB subunit near ec $3_B$ ) alters the electron transfer ratio of the two branches, suggesting that in Photosystem I, the rates are very sensitive to the energetics established by the pigment-protein interactions (Li et al., 2006).

#### B. Coupling

In addition to controlling the energies of the cofactors, the protein must provide an environment that establishes a suitable coupling, or formally the effective electronic matrix element, between the cofactors poised in their excited or charged states for the reactions. In theoretical models, coupling can be considered as arising from a pathway from the donor to the acceptor. In these models, the protein is considered to consist of a complex array of bonds as well as close but non-bonding contacts. A theoretical framework has been developed to express the protein as a network of covalent bonds, hydrogen bonds, and through-space jumps along which electron transfer proceeds in many steps (Lin et al., 2005; Gray and Winkler, 2005). Rather than making use of explicit calculations, coupling is more commonly estimated by assuming that the protein environment is relatively homogeneous, with the coupling being primarily dependent upon only the distance between the donor and acceptor (Gray and Winkler, 2005; Noy et al., 2006). In this simplified case, the coupling exponentially decreases with increasing separation (Fig. 5) showing that the maximal electron transfer rates can span a range of nearly  $10^{12}$  s<sup>-1</sup> for cofactors in close contact to a rate of 25 s<sup>-1</sup> when the separation is 23 Å.

Because of the exponential dependence on distance, electron transfer in biological systems often makes use of a series of steps involving closely positioned cofactors rather than utilizing a single step between more distantly located cofactors. In order to meet the critical demands for efficient electron transfer, careful positioning of the cofactors and fine control of the energies of the cofactors is essential. For this reason, electron transfer in the reaction center from  $P^*$  to  $Q_B$  involves a series of intermediate acceptors,  $B_A$ ,  $H_A$ , and  $Q_A$ . In general, the initial steps have very fast electron transfer rates between closely aligned cofactors. The approximate two-fold symmetry for the two branches suggests comparable couplings

as has been found by consideration of the relative rates in various mutants (Katilius et al., 2002b). Comparison of the P\* to  $P^*H_A^-$  and P\* to  $P^*H_B^-$  rates and free energy differences indicates that the slower rate for the B-branch transfer is not accounted for entirely by the energetics, and so the coupling does have a contribution, although modest (Kirmaier et al., 2001, 2005).

The charge recombination reactions illustrate how the maximal rate is primarily determined by the distance. In the wild type, the two quinones have observed recombination rates that differ by a factor of 10, although the two quinones are approximately equidistant from P. However, charge recombination from  $P^+Q_B^-$  has a significantly larger reorganization energy than that from  $P^+Q^-_{A}$ , as evident by the peak of the parabola being at a higher free energy difference, resulting in very comparable maximum rates as predicted by their similar distances to P (Allen et al., 1998).

These theoretical models to describe electron transfer within proteins can be expanded to also understand electron transfer between cofactors found in two different proteins (Lin and Beratan, 2005; Lin et al., 2005; Miyashuta et al., 2005). In interprotein transfer, additional factors must be considered, such as the docking of the proteins and the involvement of water molecules at the protein interface. For example, cytochrome  $c_2$  is a water-soluble protein that transfers an electron only after binding to the reaction center (Chapter 17, Axelrod et al.). In this case, coupling is established by the ability of interfacial water molecules to form multiple hydrogen bonding pathways that connect the tunneling pathways for each of the proteins (Miyashita et al., 2005).

#### C. Dynamics

The conventional Marcus theory (Eq. 1) rests on the assumption that the electronic transition from the initial to the final state is much faster than nuclear motion. Theories have been developed that explicitly account for dynamics (Sumi and Marcus, 1986; Warshel et al., 1989; Gehlen et al., 1994). However, dynamical contributions are rarely used because of difficulties in experimentally measuring the contributions. Although the initial electron transfer has been well characterized, including the free energy dependence for the rate (Haffa et al., 2002), certain aspects of the transfer have been puzzling. The kinetics of the changes in the spectra of the pigments are complex, and the three-fold increase in the electron transfer rate as the temperature decreases from 298 to 4 K is unusual. To determine the contribution of dynamics to the initial rate, this reaction has been analyzed by probing nearby amino acid residues, measured by absorbance changes in tryptophan residues (Wang et al., 2007). These measurements show that the protein dynamics is the limiting factor in the initial electron transfer rate rather than the energetics. In the classical picture of electron transfer, the protein undergoes structural changes only after the electron transfer to the acceptor is complete (Fig. 6). This observed dependence on the dynamics shows that the protein motions occur during the process and control the observed rate by overcoming the activation barriers described in the classical Marcus theory.

#### **IV. Pathways of Electron Transfer**

#### A. B-side Electron Transfer

In wild-type reaction centers, electron transfer proceeds along the A branch with a nearly 100% efficiency. Manipulation of the various factors influencing electron transfer comes in most dramatically in experiments aimed at getting the electron to go the 'wrong' way. Mutations that increase B-branch electron transfer have been reviewed by Wakeham and Jones (2005). Achieving a significant amount of transfer along the B branch requires a combination of mutations, such as blocking the A side by mutations that prevent incorporation of  $Q_A$  or  $H_A$ , altering the energies of the A-side states to make electron transfer less favorable by mutations near  $B_A$  and  $H_A$ , and altering the energies of the B side states to make electron transfer more favorable by mutations near B<sub>B</sub> (Fig. 7). Although reaction centers from *Rba*. *sphaeroides* and *Rba. capsulatus* have very similar electron transfer rates, the effect of many subtle differences involved in the branching is illustrated by a lower propensity of reaction centers from *Rba. sphaeroides* for electron transfer along the B branch (Kirmaier et al., 2002b). In early experiments on altering the branching ratios, only electron transfer to  $H_B$  was observed, but now transfer to  $Q_B$  has been achieved with reasonable efficiency.

The difference in the energy levels of the states  $P^+B_A^-$  and  $P^+B_B^-$  with respect to  $P^*$  is an important factor in determining the branching ratio. Mutations that only increase the lifetime of  $P^*$  are insufficient

*Fig. 6.* Light results in formation of the charge-separated state  $P<sup>+</sup>H<sub>A</sub>$  in reaction centers. In the traditional static view of electron transfer, the movement of dipoles from protein side chains and water molecules is considered to be very slow compared to the initial electron transfer rate. In contrast, in the dynamic model of the reaction center, dipoles rearrange in response to the lightinduced formation of  $P_{CT}^*$ , the excited state of P, which has a charge-transfer character, followed by transfer of the electron. Recent measurements of the dynamics of the reaction center show that these rearrangements play a critical role in determining the rate of the electron transfer process (Wang et al., 2007).

for switching to B-side electron transfer (de Boer et al., 2002). Changing the energetics along the two branches can be accomplished by the introduction of an aspartic acid residue near  $B_A$  (for Gly M201 in *Rba. capsulatus*, M203 in *Rba. sphaeroides*), showing that individual residues can influence the balance between A-side electron transfer, B-side electron transfer, and charge recombination (Heller et al., 1995). Similarly, a lysine residue introduced near  $B_B$  results in an alteration of the branching ratio (at Ser L178 in *Rba. capsulatus*) (Kirmaier et al., 1999), as does a combination of mutations that destabilizes  $P^+B_A$  and stabilizes  $P^+B_B$  (at L181 in *Rba*. *sphaeroides* and *Rba. capsulatus*, and at M210 in *Rba. sphaeroides*, M208 in *Rba. capsulatus*) (Kirmaier et al., 2004). The  $\phi_B$  mutant, in which  $B_B$  is replaced with a bacteriopheophytin (because of the His M182 to Leu mutation in *Rba. sphaeroides*), is observed to form the state  $P<sup>+</sup>B<sub>B</sub><sub>B</sub>$  with a yield of 35%, although no electron transfer to  $H<sub>B</sub>$  is observed, probably because





*Fig. 7.* In wild-type reaction centers, electron transfer occurs almost exclusively along the A branch of cofactors. In reaction centers with a set of four mutations, a significant percentage of electron transfer is observed to occur along the B side. For the YFHV mutant in *Rba. capsulatus* (Kirmaier et al., 2003), the Phe L181 to Tyr mutation lowers the energy of P<sup>+</sup>B<sub>B</sub> while three changes make electron transfer along the A branch unfavorable: the Tyr M208 to Phe mutation raises the energy of  $B_{\lambda}$ , the Leu M212 to His mutation results in H<sub>A</sub> being replaced with a bacteriochlorophyll (β), and the Trp M250 to Val mutation prevents the incorporation of Q<sub>A</sub>.

the state  $P^{\dagger}B_{\text{B}}^{-}$  is lower in energy than  $P^{\dagger}H_{\text{B}}^{-}$  (Katilius et al., 1999, 2002a,b, 2003). Thus raising the free energy of  $P^+B_A^-$  and lowering the free energy of  $P^+B_B^-$  has been key in promoting electron transfer down the B side to give  $P^{\dagger}H_B^{\dagger}$ .

Yields of  $P^{\dagger}H_{\text{B}}$  of greater than approximately 30% are accomplished by combining mutations near the monomer bacteriochlorophylls with mutants that impede electron transfer along the A branch. A number of these combination mutants incorporate the β mutation, in which  $H_A$  is replaced by a bacteriochlorophyll, resulting in a substantially higher energy for the P<sup>+</sup> $\beta$ <sup>-</sup> state compared to P<sup>+</sup>H<sub>A</sub>(Kirmaier et al., 2002a, 2003; Kee et al., 2006). An exceptionally high yield of 70%  $P<sup>+</sup>H<sub>B</sub>$  was achieved by joining together mutations that alter the energies of  $B<sub>A</sub>$  and  $B_B$  with changes derived from the  $D_{LL}$  mutant, for which electron transfer along the A branch is blocked because of the lack of  $H_A$  in that mutant (Chuang et al., 2006). On the other hand, relatively small changes in the energies of  $H_A$  and  $H_B$ , such as by addition of a hydrogen bond to  $H_B$  to lower the energy level of

 $P<sup>+</sup>H<sub>B</sub>$ , are insufficient to increase the yield of B-side transfer (Kirmaier et al., 2002a).

A number of mutants with changes near the bacteriochlorophylls and bacteriopheophytins produce  $P<sup>+</sup>H<sub>B</sub><sub>B</sub>$  but do not have significant electron transfer to  $Q_B$ , attributed to the forward reaction being unable to compete with charge recombination (de Boer et al., 2002). The rate of electron transfer from  $P^{\dagger}H_B^-$  to  $P^+Q_B^-$  is much slower than the rate from  $P^+H^-_A$  to  $P^+Q^-_{A^*}$ and so competes less well with charge recombination (Kirmaier et al., 2003; Kee et al., 2006). Several factors could influence the difference in the two rates (Kirmaier et al., 2003; Paddock et al., 2005). For example, the much larger reorganization energy associated with formation of  $Q_B^-$  compared to  $Q_A^$ as discussed above, would favor a much faster rate for  $Q_A^-$  formation. Also contributing could be the free energy differences for the bacteriopheophytin to quinone electron transfer on the two sides being slightly different. Although the bacteriopheophytin to quinone separation distances for the two sides are comparable, differences in the coupling are suggested

by distinct interactions of the quinones with nearby aromatic amino acid residues (Trp M252 in *Rba. sphaeroides* and M250 in *Rba. capsulatus* near  $Q_A$ and Phe L216 near  $Q_B$ ).

Electron transfer can get all the way to  $Q_B$  via Bbranch electron transfer if the A branch is blocked (Laible et al., 2003). Removal of  $Q_A$  also allows one to observe spectra of  $Q_B$  generated only by B-side transfer (de Boer et al., 2002; Breton et al., 2004; Paddock et al., 2005, 2006). The yield of  $Q_B^-$  when the A branch is blocked depends on what other mutations are present (Wakeham et al., 2003, 2004; Frolov et al., 2005). Mutations near  $Q_B$  that decrease the polarity (Glu L212 to Ala and Asp L213 to Ala in *Rba.*  $sphaeroides$ ) increase the yield of  $Q_B^-$  by B-branch electron transfer, postulated to be either by changing the rate through altering the oxidation-reduction potential and therefore the driving force or altering the reorganization energy (Wakeham et al., 2003).

Although coupling has been modeled extensively, its effects on establishing the branching of electron transfer are somewhat difficult to measure experimentally, as it is not clear how to rationally change the spatial orbital overlaps. However, the heterodimer mutations afford a way to change the electron density distribution, which is also involved in the coupling. Mutants constructed by Kirmaier and coworkers (2005) to assess the coupling involve combining either the A or B side heterodimers and the  $\beta$  mutation that replaces  $H_A$  with a bacteriochlorophyll. This combination of mutations has the effect of making the states  $P^+B^-_A$  and  $P^+H^-_B$  approximately equal in energy (thus equalizing the free energy contribution) while raising the monomer states such that the mechanism ( superexchange) of the initial electron transfer should be the same on both branches. Although the yield of B-side transfer is low in both cases, it is approximately four-fold less for the mutant with the A-side heterodimer. Because the energies for the heterodimer mutants with the bacteriochlorophyll on either side are assumed to be comparable, the difference in yields is attributed to a small but real difference in coupling of the dimer to  $H_A$  and  $H_B$ . This coupling difference should contribute to the A-side preference for electron transfer found in other mutants and wild type.

The pathway of electron transfer also appears to be dependent on the wavelength of the excitation energy. Blue light excitation into the Soret band gives rise to a transient  $B_B^+H_B^-$  state, which decays in picoseconds at room temperature but is longer lived at low temperature (Lin et al., 2001; Haffa et al.,

2003). Although energy transfer is fast, alternate photochemistry upon differing excitation wavelengths shows that energy transfer and electronic relaxation do not always precede electron transfer (Wang et al., 2006). These surrogate conduits for the dissipation of excitation energy suggest that the function of the B branch remains to rapidly quench higher excited states (Lin et al., 2001). Thus the B branch is poised to form charge-separated states but not to result in a terminal electron transfer as the A side is. This is consistent with the location of the carotenoid, also thought to be involved in photoprotection (Lin et al., 2003), on the B branch.

The pathway of electron transfer, in addition to being dependent on wavelength, can also be dependent upon electrostatic interactions involving a cofactor and a nearby charged amino acid residue. Introduction of protonatable residues (His L168 to Glu and Asn L170 to Asp in *Rba. sphaeroides*) results in a pH dependence for the states generated by light excitation (Haffa et al., 2004). Because the charges on the ionizable residues depend on their protonation state, and the protonation state is pH dependent, the pathway then becomes pH dependent, demonstrating an ability to switch the pathway by adjustment of the pH.

In summary, the highly efficient transfer of electrons along the A branch in wild type can be manipulated by mutagenesis to generate charge-separated states involving the B-branch cofactors. Overall, electron transfer to  $H<sub>B</sub>$  is largely limited by the unfavorable energy of  $B_B$  and its limited coupling to P<sup>\*</sup>. Electron transfer to  $Q_B$  can be observed but only when the normal electron transfer pathway along the A branch is blocked. The yield of electron transfer to  $Q_B$  is influenced by several factors, but most likely is limited by the rate being slow due to a high reorganization energy associated with formation of  $Q^-_{B}$ .

#### B. New Electron Transfer Reactions

 Whereas most of the emphasis concerning electron transfer in reaction centers has been on elucidating the factors that control the rates and the asymmetry of electron transfer, other manipulations have introduced new pathways into the reaction center (see also Chapter 16, Jones). The focus in these experiments has been on exploiting the environment near P so as to introduce novel secondary donors. In *Rba. sphaeroides*, P is normally reduced by a water-soluble cytochrome  $c_2$  (see Chapter 17, Axelrod et al.). In the absence of the cytochrome, P remains oxidized for approximately 1 s until charge recombination from  $Q_B^-$  occurs. By comparison, after light excitation and charge separation, the oxidized primary electron donor of Photosystem II is rapidly reduced by a redox active tyrosine (Tyr 161 of the D1 subunit), which is usually identified as  $Y_z$ . The electron transfer continues with  $Y^{\dagger}$  being reduced by the Mn cluster where the four equivalents are stored until oxygen is evolved. Because of the relatively low midpoint potential of 500 mV for P, the bacterial reaction centers are normally not capable of oxidizing tyrosine or manganese. However, after hydrogen bonds are added to the dimer, P becomes highly oxidizing and is capable of oxidizing tyrosine residues, including a tyrosine introduced at a location analogous to that of  $Y_z$  (Kálmán et al., 1999, 2003a,b; Narváez et al., 2002, 2004). Likewise, the modified reaction centers are capable of oxidizing manganese, which can be bound after amino acid residues acting as ligands are introduced (Fig. 8) (Thielges et al., 2005; Kálmán et al., 2005). These results demonstrate that new electron transfer pathways can be introduced, and such approaches can be used to delineate the complex evolutionary process by which anoxygenic photosynthetic complexes evolved into oxygenic complexes.

#### **V. Conclusions**

When a reaction center is excited by light, the interplay between the cofactors and the surrounding protein determines the route of electron transfer. The remarkable robustness of the reaction center to both biochemical manipulation and mutagenesis provides the opportunity to probe the features that influence electron transfer. The standard route down the A branch is observed to have the characteristics of energetics, coupling and dynamics that make it a durable electron transfer pathway. However multiple perturbation of the system makes possible other opportunistic reactions.

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*Fig. 8.* Comparison of the three-dimensional structures of the Mn binding mutant (Thielges et al., 2005) and Photosystem II shows many similar features. In both cases, three critical amino acid residues are present, a carboxylate (M173 and D1-170), a tyrosine (M164 and  $Y_z$ ), and a histidine (M193 and D1-190). In addition, both complexes have a bound redox active Mn (a mononuclear and a Mn<sub>4</sub> cluster for the mutant and Photosystem II respectively; Mn atoms shown as small spheres) that can serve as a rapid secondary electron donor to the oxidized primary donor.

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