



Crystal structure and electrical transport of nano-crystalline strontium-doped neodymium ortho-ferrites

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Abstract Strontium-doped neodymium orthoferrites belong to the perovskite type of materials of a wide range of interests and applications. This material has an orthorhombic perovskite crystal system of space group (Pbnm) no. 62. $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ is synthesized using the co-precipitation method, and the resulting precipitated powder was fired at different temperatures to control the nano-crystalline size. The effect of the nano-crystalline size on the structure and the electrical properties were studied. The EDS spectra confirmed that the compound was formed according to the suggested calculations. SEM micrographs showed the homogeneity of the nano grain size in the compound. The nano-crystalline size of the prepared samples, fired at 705 °C and at 850 °C were 40 nm and 90 nm, respectively. $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ showed a semiconducting behavior according to the exponential dependence of resistiv-

ity on temperature in the range from 350 K up to 473 K. The temperature independence of impedance is found at high frequencies starting from 0.5 MHz. The frequency dependence–activation energy is found in the temperature range from 350 K up to 473 K. The found decrease in AC resistivity with frequency increase for $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ showed that the hopping mechanism occurred in the charge carrier transfer.

Keywords Perovskites · Nano-crystalline · Neodymium · Orthoferrites · Dielectric properties

Introduction

The well-known rare earth orthoferrites, orthomanganites, orthocoblates, or orthonicklates are very rich materials because they have very interesting properties as well as a wide range of applications (Bashkirov Sh et al. 2005; Abdel-Latif et al. 2008; Hamad et al. 2017; Abdel-Latif et al. 2018a; Kurbakov et al. 2019). All of these interests were the motivation for more scientific curiosity to pursue research on these interesting topics in the hope of discovering more applications of these materials. Moreover, great efforts were carried out to find interpretations and propose new models that explain the unique and unusual phenomena found in these materials. These wide class of materials could be formed in different crystalline structures and showed crystalline phase transformations because of different annealing temperatures (Abdel-Latif et al. 2015; Kumar 2016). They are stable materials, and

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they showed very interesting chemical properties such as chemical sensing, (Shuk et al. 1993; Abdel-Latif et al. 2018b; Shuk and Guth 1995) and photocatalytic properties (Abdel-Latif et al. 2019; Abdel-Latif et al. 2017; Kanhere 2014). Both magneto-electrical and magneto-caloric properties are found for these materials such as magnetoresistance (Bouziane et al. 2005; Petrov et al. 1999), magneto-caloric effect (Ahmed et al. 2019; Bettaibi et al. 2016); moreover, multiferroics properties that are clear in these types of materials (Lee et al. 2008; Ramesh and Spaldin 2007; LEE et al. 2005). All these properties have opened the door to apply them in magnetic devices (Abdel-Latif et al. 2016; Zhang et al. 2016; Tugova et al. 2017); magnetic cooling (Ayas 2017; Chandran et al. 2019; Cherif et al. 2014) and magnetic storage media (Kundu 2016; Kim et al. 2019; Pattnaik et al. 2019). They have different crystallographic crystal systems (cubic, orthorhombic, hexagonal, monoclinic, and rhombohedral crystal system) and this is the main reason for the various characteristics of these materials (Elghoul et al. 2018, Mahato Dev et al. 2016, Tokunaga et al. 2006). Rare earth ferrites are shaped with various crystallographic symmetries (Bashkirov Sh et al. 2003; Parfenov et al. 2003; Abdel-Latif 2011; Parfenov 2007; Yousif et al. 2011; Abdel-Latif 2016; Markovich et al. 2013; Iqbal et al. 2012). Type of the formed crystallographic symmetry depends on two main parameters: the temperature treatment and the ionic radius of the A-site element and doping divalent element or rare earth element (Abdel-Latif 2020). The orthorhombic distortion of the perovskite-like structure ABO_3 is one of the common structures these compounds have. The A sites in such materials are surrounded by quite distorted twelve oxygen atoms polyhedral while the B atoms are less distorted and they are surrounded by six oxygen atoms forming an octahedron. When an octahedron in these compounds is tilted, the occurring tilt causes tilting of the neighboring octahedra. The magnetoresistance that originated because of hole doping is shown in the orthorhombic the rare earth ferrite type (so called orthoferrites).

The dielectric loss factor measurements may help to understand different types of polarization mechanisms such as Maxwell–Wagner, ionic, dipole, and electronic (Saleh et al. 2020) as well as electric conduction. Grain boundaries formed during the annealing process play a fundamental role in the crystallization and the low surface area besides the oxidation state (Iqbal et al. 2017). The grain boundary contribution is noted and clear when the high resistivity of

under-investigation materials at low frequency is measured (Naseem et al. 2018).

In the present work, $Nd_{0.6}Sr_{0.4}FeO_3$ is prepared by using the co-precipitation method and it is fired at various temperatures in order to get grains with different sizes. The effect of the nano-crystalline size on electrical and dielectric properties in $Nd_{0.6}Sr_{0.4}FeO_3$ is studied.

Experimental details

The strontium-doped neodymium orthoferrite compound under study was synthesized using the co-precipitation method. Highly pure chlorides (purity 99.9%, Sigma-Aldrich) were used: $FeCl_3 \cdot 4H_2O$, $NdCl_3 \cdot 6H_2O$, and $SrCl_2 \cdot 4H_2O$. These pure chlorides (16.68 g of $FeCl_3 \cdot 4H_2O$, 5.38 g of $NdCl_3 \cdot 6H_2O$, and 2.68 g mol of $SrCl_2 \cdot 4H_2O$) are dissolved in 100 mL distilled water to react with 5 g of NaOH. The mixed solution was kept stirred for 6 h at temperature of 85 °C, and we adjusted the pH to be 12. The precipitation process was started; the resulting precipitated powder needed separation, and so it was washed by distilled water many times, until NaCl was removed and only the formed strontium-doped neodymium orthoferrite remained. The separated powder is carefully milled to be pressed into the disc shape of 10-mm diameter. This disc is fired at the start of the annealing process at different temperatures. The firing temperatures in our case were 750 and 850 °C, in air for 12 h, and cooling was done at furnace temperature. The XRD patterns were measured using a PANalytical X'Pert Pro MPR diffractometer which produced a Cu radiation of 1.54 Å wavelength. The morphology of the compound was studied using a field emission scanning electron microscope (FE-SEM: JEOL JSM-7600F). The energy-dispersive EDS spectrum was used to examine the chemical purity of the formed $Nd_{0.6}Sr_{0.4}FeO_3$. Both the AC and DC resistivity temperature-dependent measurements were done in the temperature range from room temperature up to 473 K. The dielectric characteristics were collected as a function of temperature and frequency (LCR) meter (HP4284A). The DC resistivity measurements were carried out as a function of temperature using the two- and four-point probe technique (Scientific Equipment Service).

Results and discussions

The crystallization of rare earth ferrite, manganite, cobaltite, or nickelate oxides occurred in a certain way as to take the perovskite-like structure (Bashkirov Sh et al. 2005; Kurbakov et al. 2019; Abdel-Latif et al. 2018b; Abdel-Latif et al. 2019). This type of perovskite-like structure can be found in different crystal systems where it takes the simple cubic form and transforms into another crystal structure phase, such as orthorhombic, hexagonal, monoclinic, or rhombohedral structure. This difference is attributed to the method of synthesis and the annealing process during preparation. The great diversity in crystallization in various forms leads to a wide scale of interests correlated to each structure form as well as the change in the crystal structure at low temperature. The EDS test of the strontium-doped neodymium orthoferrite compound under study showed that it is formed according to the calculated formula and all the constituent elements in the compound are formed with the right percentage.

XRD structure analysis

In the present work, XRD patterns of $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ are measured at room temperature for as-prepared samples as well as samples fired at 750 °C and at 850 °C. These measurements were done using a Cu X-ray tube of wavelength 1.54 Å according to a step scan regime from angle $2\theta = 10^\circ$ up to 90° . The XRD patterns were analyzed to check the forming of the designed orthorhombic phase structure. It is clear from Fig. 1a that the orthorhombic phase is formed after annealing at 750 °C and 850 °C temperatures. All the peaks in the XRD patterns are indexed, and it is clear that the orthorhombic crystal structure system of space group no. 62 (Pbnm) is satisfied.

The crystalline phase is not completely formed in the as-prepared sample, but because of the annealing process (at 750 °C and 850 °C), the crystalline phase is completely formed. The formation required in the crystal structure as a result of annealing is observed in the nanoscale. The experimental XRD patterns of the fired samples are fitted based on the Rietveld refinement method by using FullProf package software (Rodriguez-Carvajal 1993), and all the calculated lattice parameters from the refinements are listed in Table 1. The indexing of all the reflection peaks in the measured XRD patterns is done according to orthorhombic crystal

system of space group Pbnm (no. 62). All generated hkl planes are shown in the CIF file as well as the Rietveld refinement parameters.

The resulting lattice constants of $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ fired at 750 °C and 850 °C, which are listed in Table 1, are in good agreement with those previously reported (Dasgupta et al. 2002). There is slight change in the lattice constant, bond lengths, and angle between iron and oxygen atoms. This difference is attributed to the difference in the annealing temperature which may affect how oxygen atoms are occupying their positions and forming the octahedral side around the iron atoms. There are differences in the bond lengths and angle between iron and oxygen which are the two main parameters in electronic and magnetic transfer. On the other hand, the difference in crystalline size in both samples is clear because of the various annealing temperatures. The crystalline size could be deduced from the well-known Scherrer equation where the broadness of the peaks is analyzed and taking into consideration the instrumental broadening (Abdel-Latif 2011) to avoid a dramatical error:

$$\text{Cryst size } \delta = \frac{K\lambda}{B \sin\theta} \quad (1)$$

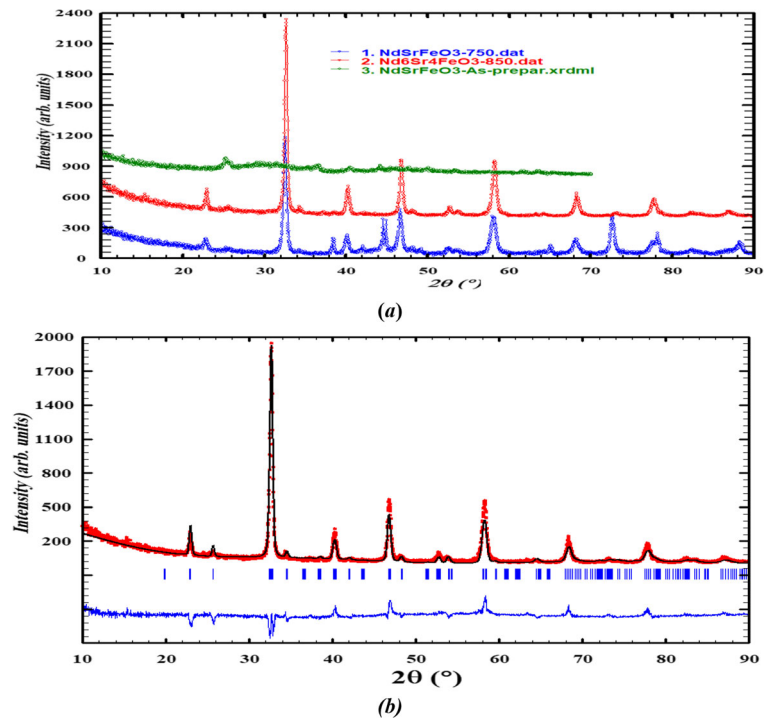
where the broadening in the peak (FWHM) is represented by B which is equal to the difference between that measured in the experimental XRD pattern for sample B_{obs} and that measured for the standard sample B_{std} . The importance of this crystalline size is because it plays a principal role in the dependence of the physical and chemical properties.

Electron density $\rho(r)$ calculations as a result of the scattering of incident X-ray on the unit cell are done using a subprogram in FullProf software package. These calculations are based on Fourier method and fast Fourier transform (FFT). The following equation was applied for the density of electrons (DOE) based on FFT in the fitted XRD data, estimating the density of electrons (Abdel-Latif et al. 2016):

$$\rho(r) = \frac{1}{V} \sum_H F(H) \exp\{-2\pi i(H \cdot r)\} \quad (2)$$

It is clear that the electron density equation consists of the following parts: for the unit cell, there are the volume (V), the position vector inside the unit cell (r), and the reciprocal lattice vector (H) besides the coefficients of the complex Fourier $F(H)$. The density of

Fig. 1 XRD patterns of **a** as-prepared samples as well as samples fired at 750 °C and at 850 °C at room temperature and of **b** $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ fired at 850 °C with Rietveld analysis at room temperature



electrons has the same units of the complex Fourier $F(H)$ per volume (Abdel-Latif 2020). From the distribution of electrons inside the unit cell, it is clear that there are more electrons in the position of iron atoms in the case of annealing at 850 °C than in the case of annealing at 750 °C. This difference in the electron density will help us to understand the electric and magnetic properties dependent on the annealing temperature (Fig. 2).

Table 1 Lattice parameters, crystalline size, bond lengths, and angles between iron and oxygen in the $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ samples fired at 750 °C and at 850 °C

| $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ | 750 °C | 850 °C |
|--|-------------|-------------|
| a (Å) | 5.483 (4) | 5.4709 (8) |
| b (Å) | 5.536 (4) | 5.5158 (6) |
| c (Å) | 7.748 (6) | 7.7480 (12) |
| V (Å ³) | 235.2 (3) | 233.81 (6) |
| Cryst size (nm) | 24 | 72 |
| Fe-O ₂ (Å) | 1.957 (19) | 2.02 (3) |
| Fe-O ₂ (Å) | 1.999 (19) | 2.06 (3) |
| Fe-O ₁ (Å) | 1.9470 (15) | 1.9469 (3) |
| Fe-O ₂ -Fe | 160.0 (8) | 145.1 (11) |
| Fe-O ₁ -Fe | 168.39 (19) | 168.42 (4) |

The representation of the unit cell based on the CIF file which resulted from Rietveld refinement is done using Diamond software, and it is shown in Fig. 3. It is clear that there is a slight difference in the atomic position in the unit cell which may effect its properties.

EDXS and SEM analysis

The elemental analysis based on EDXS spectra was done for $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ as shown in Fig. 4. The used accelerating voltage for collecting spectrum was 15 keV. Analyzing the intensity under the peaks for all elements in the spectrum, one found that there is good agreement between the experimentally observed percentages of elements and the theoretically calculated. This good agreement showed that the obtained compound is formed in the proper way according to the proposed structure and all elements existed at the right percentage. The nano-crystalline size from the SEM micrographs of $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ annealed at 750 °C and 850 °C is in good agreement with those calculated from XRD analysis in the range from 20 nm up to 100 nm. The average size is 40 nm for NSFO-750 °C and 90 nm for NSFO-850 °C.

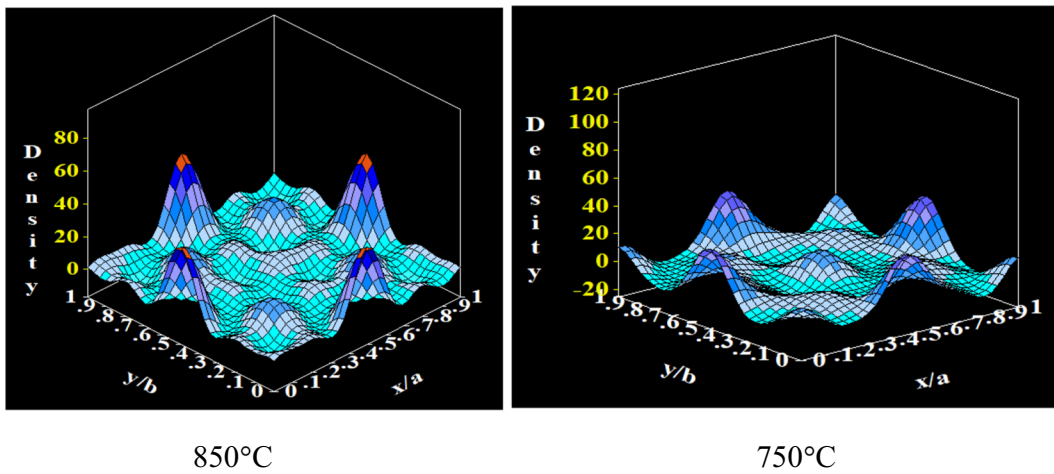


Fig. 2 Electron density map representation inside the unit cell of $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ fired at 750 °C and 850 °C

Transport properties of $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$

To understand the electrical transport in the strontium-doped neodymium orthoferrite resistivity under investigation, measurements for temperature dependence were done. The variation in AC resistivity measurements as a function of temperature are shown in Figs. 5, 6, 7, and 8.

The exponential relation between resistivity and temperature is found in the temperature range from 350 K up to 473 K to show the semiconducting properties of these materials under study (Parfenov 2007; Yousif et al. 2011).

The AC resistivity as a function of frequency and temperature was analyzed for $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$. The

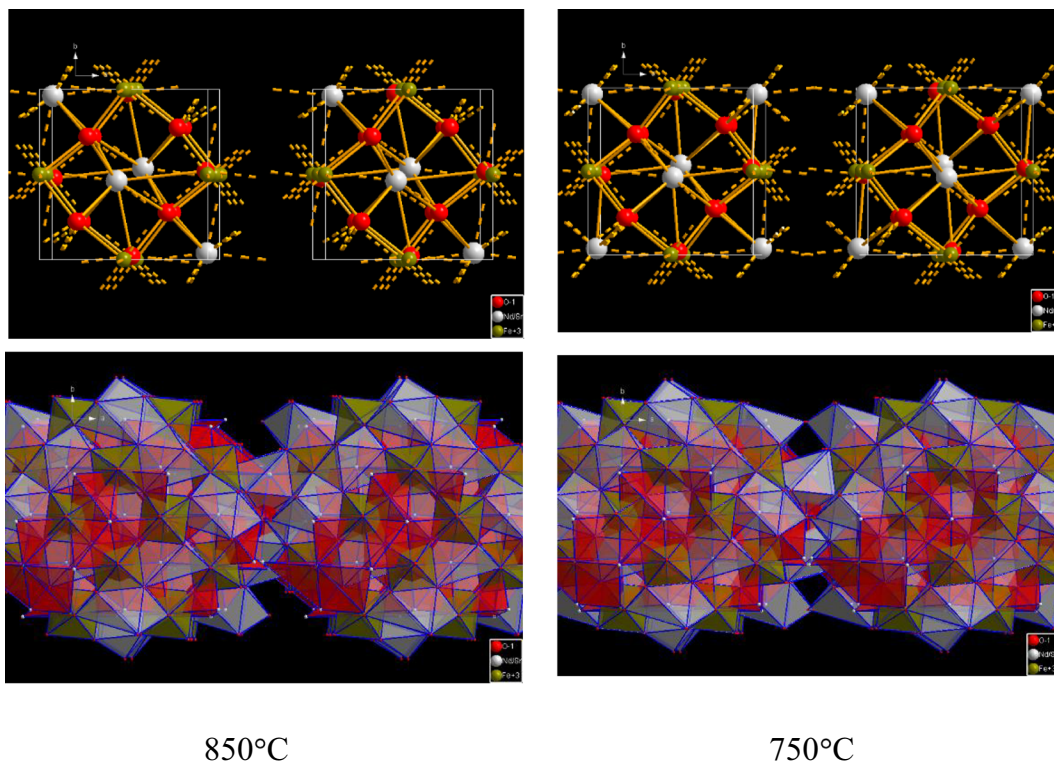


Fig. 3 Unit cell representation of $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ fired at 750 and 850 °C

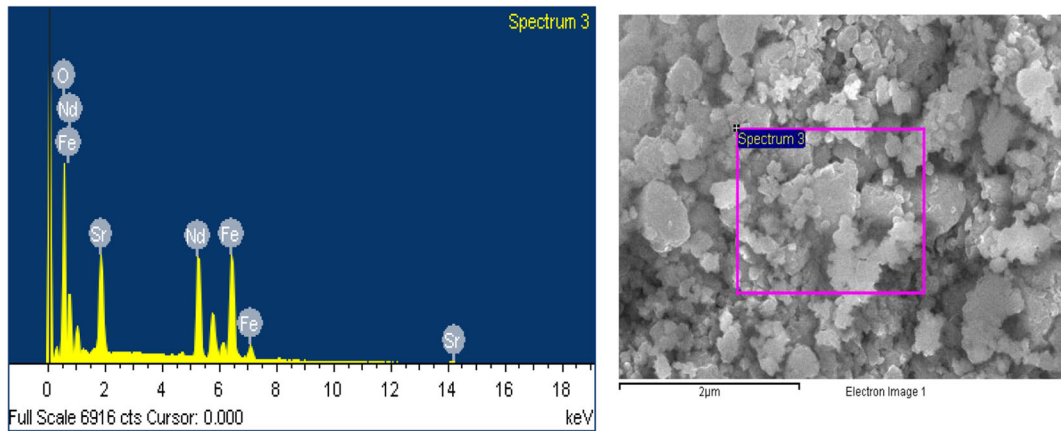


Fig. 4 EDS spectrum and SEM micrograph of $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$

decrease in AC resistivity in NSFO-750 °C as a result of the increase in frequency is clear as shown in Fig. 6. The present behavior enhances the hopping mechanism in NSFO-750 °C (Rajwali and Fang 2015). This type of AC conductivity–frequency dependency may be interpreted in terms of series resistance effect according to what was reported before. The strong correlation between AC resistivity of any materials and the hopping of weakly bound charge carriers is clear in the NSFO-750 °C under investigation and is similar to those reported in Ahmad et al. (2018).

The frequency dependence–activation energy is represented in Fig. 8. The obtained results confirm the semiconducting behavior of NSFO-750 °C in the temperature range from 350 K up to 473 K. The activation

energy at low frequency is closed to each other and near the value of DC resistivity dependent on temperature (zero frequency case).

The relation between the complex impedance Z and microstructure of polycrystalline is strongly pronounced where the grain size is a very important parameter which may affect Z values. There are other parameters that may affect on complex impedance such as the temperature, frequency, and conductivity type. Knowing the conduction mechanism dependency on the complex impedance measurements helps us to understand the electronic transfer. In general, the electrical properties may be expressed in terms of two parameters: the real and imaginary parts of impedance (Dadami et al. 2017)

Fig. 5 AC resistivity variation as function of temperature for different frequencies

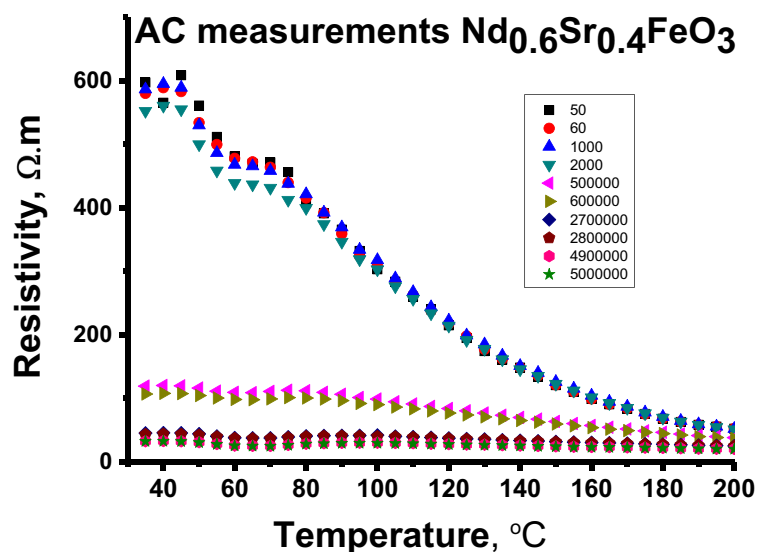
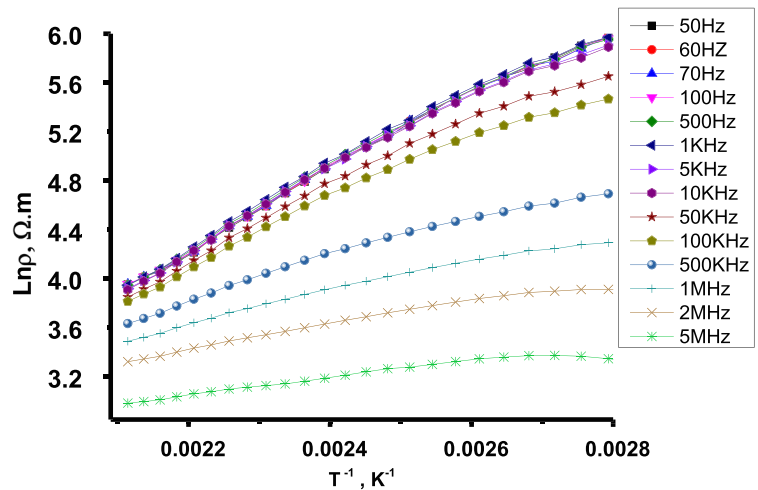


Fig. 6 The linear dependence of $\ln \rho_{ac}$ on inverse temperature for NSFO-750C at different frequencies



$$Z^* = Z' + jZ'' \tag{3}$$

where the real part is defined as $Z' = |Z| \cos \theta$ and $Z'' = |Z| \sin \theta$ is the imaginary part. Impedance-temperature dependence measurement at various frequencies of $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ annealed at 750°C (NSFO-750 $^\circ\text{C}$) is shown in Fig. 5.

The high value of the impedance at low frequency is clear and followed by a sharp decrease when both temperature and frequencies increased. In other words, we can note the temperature independence of impedance which has been started from 0.5 MHz at high frequencies.

Electric conduction and different polarization mechanisms are parameters that depend on the dielectric loss factor (Komarov et al. 2005). Permittivity is a physical parameter that is used to find the

dielectric properties. From the electromagnetic wave reflection at interfaces and the occurring attenuation of wave energy within materials, we can get permittivity and may be expressed as

$$\epsilon^* = \epsilon' + j\epsilon'' \tag{4}$$

where ϵ' is the real permittivity and ϵ'' is the imaginary permittivity in the equation of permittivity (Saleh et al. 2020).

The ϵ' real and ϵ'' imaginary parts of the dielectric dependence on frequency can be written using the impedance data accordingly:

$$\epsilon' = \frac{Z''}{2\pi f C_0 [(Z')^2 + (Z'')^2]} \tag{5}$$

$$\epsilon'' = \frac{Z'}{2\pi f C_0 [(Z')^2 + (Z'')^2]}$$

where Z' is the real impedance part and Z'' is the imaginary impedance part; C_0 is the geometrical capacitance and f is the frequency.

Looking at variation in dielectric loss dependent on frequency at different temperatures shown in Fig. 9, one can see the high value of dielectric loss at very low frequencies then a sharp decrease in dielectric loss. This behavior is repeated, but the decrease is not so sharp like in room temperature. The sharp decrease in dielectric loss, with frequency increase, is less at high temperature measurements, 423 and 473 K, as clearly shown in Fig. 9, where the dielectric loss is independent on frequency at very high frequency (dielectric dispersion). This behavior can

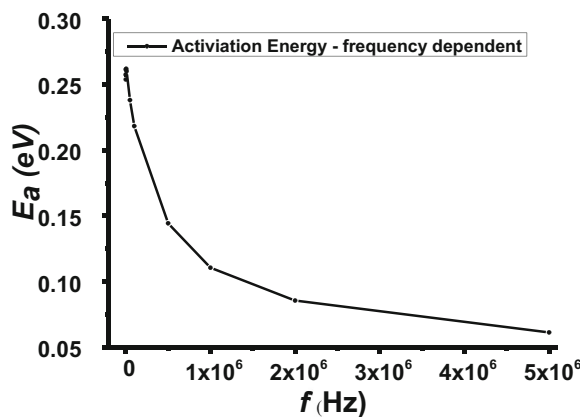
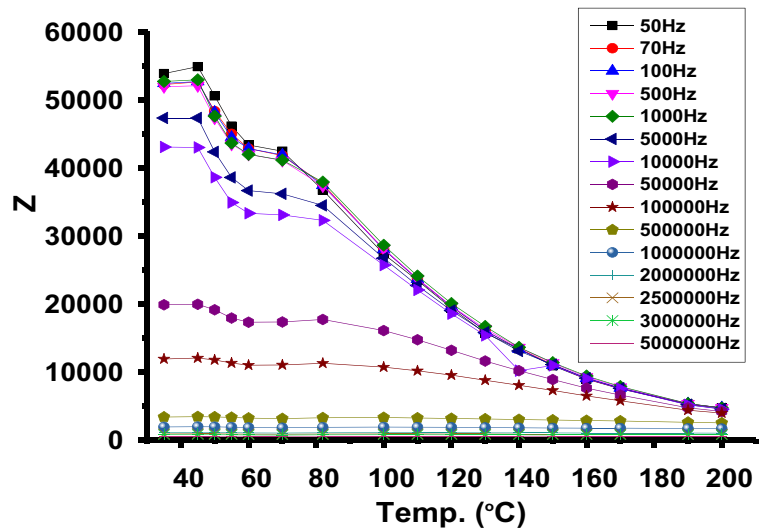


Fig. 7 Frequency dependence of activation energy E_a in NSFO-750C

Fig. 8 The correlation between the impedance and temperature of NSFO-750C at different frequencies



be interpreted in terms of conducting grains generated in nano-structure in consistent with the Maxwell and Wagner’s bilayer model and Koops phenomenological theory (Bhasin et al. 2018). Moreover, the

dielectric loss tangent $\tan \delta$ gives us information about the energy loss. It is the real permittivity divided by the imaginary dielectric permittivity equation:

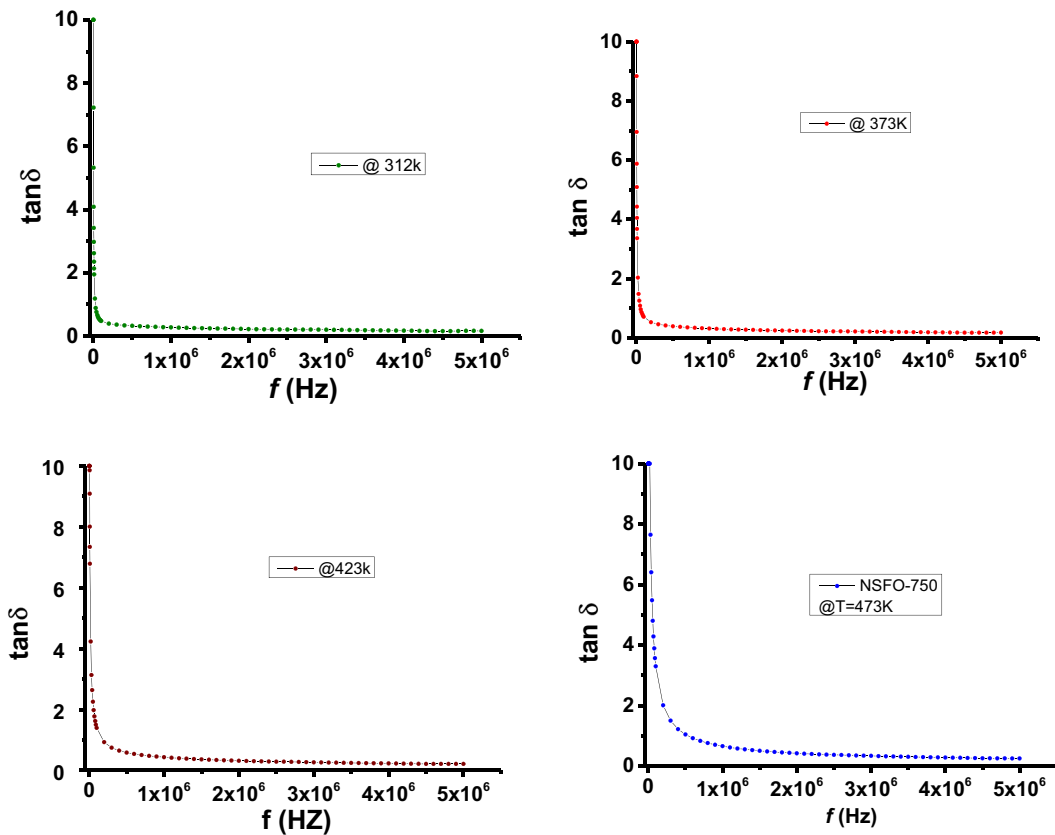


Fig. 9 Variation in dielectric loss dependent on frequency at different temperatures

$$D = \tan\delta = \frac{\epsilon''}{\epsilon'} \quad (6)$$

where δ is the phase difference between the applied field and the induced current in the material (Varshney et al. 2018). The observed decrease in $\tan \delta$, in our case, with frequency increase, is mainly based on two parameters: resistive loss and relaxation loss. In resistive loss, the mobility of charge has a great effect, but in relaxation loss, the dipole effect is more pronounced (Sharma et al. 2014).

Conclusions

- The orthorhombic crystal structure system of space group Pbnm is found.
- The obtained nano-crystalline sizes of the prepared $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ samples, fired at 705 °C and at 850 °C, were 40 nm and 90 nm, respectively.
- The found decrease in AC resistivity with frequency increase gives the interpretation for the charge carriers transfer mechanism in $\text{Nd}_{0.6}\text{Sr}_{0.4}\text{FeO}_3$ as hopping mechanism.
- The temperature independence of impedance is found at high frequencies starting from 0.5 MHz.

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Compliance with ethical standards

Conflict of interest The author declares that he has no conflict of interest.

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